

# THE HIRZEBRUCH $\chi_y$ GENUS AND A THEOREM OF HIRZEBRUCH ON ALMOST COMPLEX MANIFOLDS

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ABSTRACT. The purpose of this note is to give an introduction to the Hirzebruch  $\chi_y$  genus and to give a proof of a theorem of Hirzebruch which states that on a closed almost complex manifold  $M$  of dimension  $4m$  we have  $\chi(M) \equiv (-1)^m \sigma(M) \pmod{4}$ .

Let  $(M, g)$  be an  $2n$ -dimensional closed Riemannian manifold. Given a  $\text{spin}^c$  structure, one can form the complex  $\text{spin}^c$  bundles  $\mathbb{S}_\mathbb{C}^+$  and  $\mathbb{S}_\mathbb{C}^-$ . Then there is a  $\text{spin}^c$  Dirac operator  $\not{D}^c : \Gamma(\mathbb{S}_\mathbb{C}^+) \rightarrow \Gamma(\mathbb{S}_\mathbb{C}^-)$  which has index

$$\text{ind}(\not{D}^c) = \int_M \exp(c_1(L)/2) \hat{A}(TM)$$

where  $L$  is the complex line bundle associated to the  $\text{spin}^c$  structure; see Theorem D.15 of [5].

If  $E \rightarrow M$  is a hermitian vector bundle, then there is a twisted  $\text{spin}^c$  Dirac operator  $\not{D}_E^c : \Gamma(\mathbb{S}_\mathbb{C}^+ \otimes E) \rightarrow \Gamma(\mathbb{S}_\mathbb{C}^- \otimes E)$  which has index

$$\text{ind}(\not{D}_E^c) = \int_M \exp(c_1(L)/2) \text{ch}(E) \hat{A}(TM).$$

I don't know a reference for this precise statement (if you do, please let me know), but the fact that this quantity is an integer is Theorem 26.1.1 of [3].

Suppose now that  $M$  admits an almost complex structure and  $g$  is hermitian. Then there is a canonical  $\text{spin}^c$  structure which has associated line bundle  $L = \det_\mathbb{C}(TM)$ , so  $c_1(L) = c_1(M)$ ; see Example D.6 of [5]. Using the fact that  $\exp(c_1(M)/2) \hat{A}(TM) = \text{Td}(TM)$ , the index becomes

$$\text{ind}(\not{D}_E^c) = \int_M \text{ch}(E) \text{Td}(TM). \tag{1}$$

In addition, the complex  $\text{spin}^c$  bundles take the form  $\mathbb{S}_\mathbb{C}^+ \cong \bigwedge^{0,\text{even}} M$  and  $\mathbb{S}_\mathbb{C}^- = \bigwedge^{0,\text{odd}} M$ ; see corollary 3.4.6 of [6]. If  $E = \bigwedge^{p,0} M$ , then we have a twisted  $\text{spin}^c$  Dirac operator  $\not{D}_{\bigwedge^{p,0} M}^c : \Gamma(\bigwedge^{p,\text{even}} M) \rightarrow \Gamma(\bigwedge^{p,\text{odd}} M)$ ; for notational convenience, we will instead write  $\not{D}_p^c$  for this operator. We define  $\chi^p(M) := \text{ind}(\not{D}_p^c)$ ; if  $p = 0$ , this is just the Todd genus. The Hirzebruch  $\chi_y$  genus is defined to be

$$\chi_y(M) := \sum_{p=0}^n \chi^p(M) y^p.$$

## INTEGRABLE CASE

Suppose now that  $J$  is integrable, in which case  $n = \dim_\mathbb{C} M$ . Then, modulo order zero terms, we have  $\not{D}^c = \sqrt{2}(\bar{\partial} + \bar{\partial}^*)$ ; see Proposition 8 of [1]. In addition, if  $E$  is holomorphic, then modulo order zero terms  $\not{D}_E^c = \sqrt{2}(\bar{\partial}_E + \bar{\partial}_E^*)$  and (1) becomes the statement of the Hirzebruch-Riemann-Roch theorem.

In particular,  $\bigwedge^{p,0} M$  is holomorphic and  $\bar{\partial}_p^c : \Gamma(\bigwedge^{p,\text{even}} M) \rightarrow \Gamma(\bigwedge^{p,\text{odd}} M)$  is just  $\sqrt{2}(\bar{\partial} + \bar{\partial}^*)$  to highest order. If  $\mathcal{H}_{\bar{\partial}}^{p,q}(M)$  denotes the  $\bar{\partial}$ -harmonic  $(p,q)$ -forms on  $M$ , then

$$\begin{aligned} \text{ind}(\bar{\partial}_p^c) &= \text{ind} \sqrt{2}(\bar{\partial} + \bar{\partial}^*) \\ &= \dim \left( \bigoplus_{q \text{ even}} \mathcal{H}_{\bar{\partial}}^{p,q}(M) \right) - \dim \left( \bigoplus_{q \text{ odd}} \mathcal{H}_{\bar{\partial}}^{p,q}(M) \right) \\ &= \sum_{q \text{ even}} \dim \mathcal{H}_{\bar{\partial}}^{p,q}(M) - \sum_{q \text{ odd}} \dim \mathcal{H}_{\bar{\partial}}^{p,q}(M) \\ &= \sum_{q \text{ even}} h^{p,q}(M) - \sum_{q \text{ odd}} h^{p,q}(M) \\ &= \sum_{q=0}^n (-1)^q h^{p,q}(M) \\ &= \chi(M, \Omega^p). \end{aligned}$$

Using the penultimate expression above, we have

$$\chi_y(M) = \sum_{p=0}^n \sum_{q=0}^n (-1)^q h^{p,q}(M) y^p.$$

We now list some properties of  $\chi_y(M)$  in the integrable case.

**Property 1.**  $\chi_y(M) = (-y)^n \chi_{y^{-1}}(M)$ .

*Proof.* As

$$(-y)^n \chi_{y^{-1}}(M) = (-y)^n \sum_{p=0}^n \chi^p(M) y^{-p} = \sum_{p=0}^n (-1)^n \chi^p(M) y^{n-p},$$

this property is equivalent to  $\chi^p(M) = (-1)^n \chi^{n-p}(M)$ .

By Serre duality we have  $h^{p,q}(M) = h^{n-p,n-q}(M)$ , so

$$\begin{aligned} \chi^p(M) &= \sum_{q=0}^n (-1)^q h^{p,q}(M) \\ &= \sum_{q=0}^n (-1)^q h^{n-p,n-q}(M) \\ &= (-1)^n \sum_{q=0}^n (-1)^{n-q} h^{n-p,n-q}(M) \\ &= (-1)^n \chi^{n-p}(M). \end{aligned}$$

□

**Property 2.** *If  $M$  admits a Kähler metric, then  $\chi_{-1}(M) = \chi(M)$ .*

*Proof.* Note that in the Kähler case

$$\chi_{-1}(M) = \sum_{p=0}^n \sum_{q=0}^n (-1)^{p+q} h^{p,q}(M) = \sum_{k=0}^{2n} (-1)^k \sum_{p+q=k} h^{p,q}(M) = \sum_{k=0}^{2n} (-1)^k b_k(M) = \chi(M).$$

□

**Property 3.** *Suppose that  $n$  is even and  $M$  admits a Kähler metric. Then  $\chi_1(M) = \sigma(M)$ .*

*Proof.* Using the fact that  $h^{p,q}(M) = h^{q,p}(M)$ , we have

$$\chi_1(M) = \sum_{p=0}^n \sum_{q=0}^n (-1)^q h^{p,q}(M) = \sum_{p=0}^n \sum_{q=0}^n (-1)^q h^{q,p}(M) = \sum_{p=0}^n \sum_{q=0}^n (-1)^p h^{p,q}(M).$$

It follows from the Hard Lefschetz Theorem that the final expression is equal to  $\sigma(M)$ ; see Corollary 3.3.18 of [4].  $\square$

As we will see in the next section, all three of these properties hold in general.

### NON-INTEGRABLE CASE

Suppose now that  $J$  is not integrable.

In order to establish the properties mentioned in the previous section, we need the following expression for  $\chi_y(M)$ .

**Theorem.** *Let  $x_i$  be the Chern roots of  $TM$ . Then*

$$\chi_y(M) = \int_M \prod_{i=1}^n \frac{x_i(1 + ye^{-x_i})}{1 - e^{-x_i}}.$$

*Proof.* By the splitting principle, we can suppose that  $TM = \ell_1 \oplus \cdots \oplus \ell_n$ , and hence  $T^*M = \ell_1^* \oplus \cdots \oplus \ell_n^*$ , without any loss of generality. Defining  $x_i = c_1(\ell_i)$ , we have  $-x_i = c_1(\ell_i^*)$ . Note that

$$\bigwedge^{p,0} M = \bigwedge^p T^*M = \bigwedge^p (\ell_1^* \oplus \cdots \oplus \ell_n^*) = s_p(\ell_1^*, \dots, \ell_n^*)$$

where  $s_p$  is the  $p^{\text{th}}$  elementary symmetric polynomial (addition and multiplication correspond to direct sum and tensor product respectively). Therefore

$$\text{ch} \left( \bigwedge^{p,0} M \right) = \text{ch}(s_p(\ell_1^*, \dots, \ell_n^*)) = s_p(\text{ch}(\ell_1^*), \dots, \text{ch}(\ell_n^*)) = s_p(e^{-x_1}, \dots, e^{-x_n}).$$

So we have

$$\begin{aligned} \chi_y(M) &= \sum_{p=0}^n \chi^p(M) y^p \\ &= \sum_{p=0}^n \text{ind}(\not\partial_p^c) y^p \\ &= \sum_{p=0}^n \left( \int_M \text{ch} \left( \bigwedge^{p,0} M \right) \text{Td}(M) \right) y^p \\ &= \int_M \left( \sum_{p=0}^n \text{ch} \left( \bigwedge^{p,0} M \right) y^p \right) \text{Td}(M) \\ &= \int_M \left( \sum_{p=0}^n s_p(e^{-x_1}, \dots, e^{-x_n}) y^p \right) \prod_{i=1}^n \frac{x_i}{1 - e^{-x_i}} \\ &= \int_M \prod_{i=1}^n (1 + e^{-x_i} y) \prod_{i=1}^n \frac{x_i}{1 - e^{-x_i}} \\ &= \int_M \prod_{i=1}^n \frac{x_i(1 + e^{-x_i} y)}{1 - e^{-x_i}}. \end{aligned}$$

□

Setting  $y = -1$ , we now see that property 2 holds in the non-integrable case:

$$\chi_{-1}(M) = \int_M \prod_{i=1}^n \frac{x_i(1 - e^{-x_i})}{1 - e^{-x_i}} = \int_M \prod_{i=1}^n x_i = \int_M c_n(M) = \int_M e(M) = \chi(M).$$

When  $J$  is integrable, property 2 gives us the following result (which also follows from the existence of the Frölicher spectral sequence).

**Corollary.** *Let  $M$  be an  $n$ -dimensional compact complex manifold. The Euler characteristic of  $M$  is given by*

$$\chi(M) = \sum_{p=0}^n \sum_{q=0}^n (-1)^{p+q} h^{p,q}(M).$$

For the other two properties, we need the following lemma. Thanks to Professor Ping Li for pointing this out to me.

**Lemma.** *Let  $t$  be a parameter. Then*

$$\chi_y(M) = \int_M \prod_{i=1}^n \frac{x_i(1 + ye^{-tx_i})}{1 - e^{-tx_i}}.$$

*Proof.* The key is to note that  $\chi_y(M)$  only depends on the degree  $2n$  part of the integrand. As  $\deg x_i = 2$ , if we replace  $x_i$  by  $tx_i$ , then we have

$$\int_M \prod_{i=1}^n \frac{tx_i(1 + ye^{-tx_i})}{1 - e^{-tx_i}} = t^n \int_M \prod_{i=1}^n \frac{x_i(1 + ye^{-x_i})}{1 - e^{-x_i}} = t^n \chi_y(M).$$

Dividing through by  $t^n$ , we arrive at the result. □

With this lemma in hand, we have

$$\begin{aligned} \chi_y(M) &= \int_M \prod_{i=1}^n \frac{x_i(1 + ye^{-x_i})}{1 - e^{-x_i}} \\ &= \int_M \prod_{i=1}^n \frac{x_i(1 + ye^{x_i})}{1 - e^{x_i}} \\ &= \int_M \prod_{i=1}^n (-y) \frac{x_i(y^{-1} + e^{x_i})}{e^{x_i} - 1} \\ &= (-y)^n \int_M \prod_{i=1}^n \frac{x_i(e^{x_i} + y^{-1})}{e^{x_i} - 1} \\ &= (-y)^n \int_M \prod_{i=1}^n \frac{x_i(1 + y^{-1}e^{-x_i})}{1 - e^{-x_i}} \\ &= (-y)^n \chi_{y^{-1}}(M) \end{aligned}$$

where the second equality uses the lemma with  $t = -1$ . So we see that property 1 holds.

For property 3, suppose  $n = 2m$ . Setting  $y = 1$  gives

$$\chi_1(M) = \int_M \prod_{i=1}^n \frac{x_i(1 + e^{-x_i})}{1 - e^{-x_i}} = \int_M \prod_{i=1}^n \frac{x_i(1 + e^{2x_i})}{1 - e^{2x_i}} = \int_M \prod_{i=1}^n \frac{x_i}{\tanh(x_i)}$$

where the lemma was used with  $t = -2$  in the second equality. Recall that the power series which generates the  $L$  genus (in terms of Pontryagin classes) is

$$Q(x) = \frac{\sqrt{x}}{\tanh(\sqrt{x})}.$$

By lemma 1.3.1 of [3], the corresponding power series which generates the  $L$  genus (in terms of Chern classes) is

$$\tilde{Q}(x) = \frac{x}{\tanh(x)}.$$

So the above computation shows that

$$\chi_1(M) = \int_M \prod_{i=1}^n \frac{x_i}{\tanh(x_i)} = \int_M L_m(p_1, \dots, p_m) = \sigma(M).$$

In the integrable case, we immediately obtain the following corollary.

**Corollary.** *Let  $M$  be an  $n$ -dimensional compact complex manifold with  $n$  even. The signature of  $M$  is given by*

$$\sigma(M) = \sum_{p=0}^n \sum_{q=0}^n (-1)^q h^{p,q}(M).$$

Note, in the proof of property 3 in the Kähler case, the exponent of  $-1$  is  $p$ , not  $q$ . Because  $h^{p,q}(M) = h^{q,p}(M)$ , it doesn't make any difference. However, in the non-Kähler case, having exponent  $p$  does not compute the signature as can easily be checked for a Hopf surface.

If  $M$  has odd dimension, it follows from Serre duality that  $\chi_1(M) = 0$ .

#### A THEOREM OF HIRZEBRUCH

The three properties of  $\chi_y(M)$  give rise to the following result of Hirzebruch.

**Theorem.** *Suppose  $M$  is a closed  $4m$ -dimensional manifold which admits an almost complex structure, then  $\chi(M) \equiv (-1)^m \sigma(M) \pmod{4}$ .*

*Proof.* We separate the proof into two cases based on the parity of  $m$ .

If  $m = 2k$  is even, then  $n = 2m = 4k$  so

$$\begin{aligned} \chi(M) &= \chi_{-1}(M) \\ &= \sum_{p=0}^{4k} (-1)^p \chi^p(M) \\ &= \sum_{p=0}^{4k} \chi^p(M) - 2 \sum_{p=0}^{2k-1} \chi^{2p+1}(M) \\ &= \chi_1(M) - 2 \left[ \sum_{p=0}^{k-1} \chi^{2p+1}(M) + \sum_{p=k}^{2k-1} \chi^{2p+1}(M) \right] \\ &= \sigma(M) - 2 \left[ \sum_{p=0}^{k-1} \chi^{2p+1}(M) + \sum_{p=k}^{2k-1} (-1)^{4k} \chi^{4k-(2p+1)}(M) \right] \\ &= \sigma(M) - 2 \left[ \sum_{p=0}^{k-1} \chi^{2p+1}(M) + \sum_{p=k}^{2k-1} \chi^{2(2k-1-p)+1}(M) \right] \end{aligned}$$

$$\begin{aligned}
&= \sigma(M) - 2 \left[ \sum_{p=0}^{k-1} \chi^{2p+1}(M) + \sum_{p=0}^{k-1} \chi^{2p+1}(M) \right] \\
&= \sigma(M) - 4 \sum_{p=0}^{k-1} \chi^{2p+1}(M).
\end{aligned}$$

Therefore  $\chi(M) \equiv \sigma(M) \pmod{4}$ .

If  $m = 2k + 1$  is odd, then  $n = 2m = 4k + 2$  so

$$\begin{aligned}
\chi(M) &= \chi_{-1}(M) \\
&= \sum_{p=0}^{4k+2} (-1)^p \chi^p(M) \\
&= - \sum_{p=0}^{4k+2} \chi^p(M) + 2 \sum_{p=0}^{2k+1} \chi^{2p}(M) \\
&= -\chi_1(M) + 2 \left[ \sum_{p=0}^k \chi^{2p}(M) + \sum_{p=k+1}^{2k+1} \chi^{2p}(M) \right] \\
&= -\sigma(M) + 2 \left[ \sum_{p=0}^k \chi^{2p}(M) + \sum_{p=k+1}^{2k+1} (-1)^{4k+2} \chi^{4k+2-2p}(M) \right] \\
&= -\sigma(M) + 2 \left[ \sum_{p=0}^k \chi^{2p}(M) + \sum_{p=k+1}^{2k+1} \chi^{2(2k+1-p)}(M) \right] \\
&= -\sigma(M) + 2 \left[ \sum_{p=0}^k \chi^{2p}(M) + \sum_{p=0}^k \chi^{2p}(M) \right] \\
&= -\sigma(M) + 4 \sum_{p=0}^k \chi^{2p}(M).
\end{aligned}$$

Therefore  $\chi(M) \equiv -\sigma(M) \pmod{4}$ . □

We end with an application of this theorem.

Let  $M_k = k\mathbb{C}\mathbb{P}^{2m}$  denote the connected sum of  $k \geq 0$  copies of  $\mathbb{C}\mathbb{P}^{2m}$ . Note that  $\chi(M_k) = (2m-1)k + 2 = 2mk - k + 2$  and  $\sigma(M_k) = k$ . For  $k$  even, we have  $\chi(M_k) \equiv -k + 2 \pmod{4}$  while  $(-1)^m \sigma(M_k) \equiv (-1)^m k \pmod{4}$ . Therefore, the manifolds  $M_k$  with  $k$  even do not admit almost complex structures. On the other hand, for  $k$  odd, we have  $\chi(M_k) \equiv 2m - k + 2 \pmod{4}$  and  $(-1)^m \sigma(M_k) \equiv (-1)^m k \pmod{4}$ . By splitting into cases (either  $m$  even and  $m$  odd, or  $k \equiv 1 \pmod{4}$  and  $k \equiv 3 \pmod{4}$ ), one can verify that  $\chi(M_k) \equiv (-1)^m \sigma(M_k) \pmod{4}$  for  $k$  odd; that is, we can't use the above theorem to rule out the existence of almost complex structures on these manifolds. In fact, these manifolds do admit almost complex structures, see [2].

On the other hand, the manifolds  $N_k = k\mathbb{C}\mathbb{P}^{2m+1}$  admit almost complex structures (even integrable ones) for all  $k > 0$ . To see this, note that the map  $\mathbb{C}\mathbb{P}^{2m+1} \rightarrow \mathbb{C}\mathbb{P}^{2m+1}$  given by  $[z_0, \dots, z_{2m+1}] \mapsto [\bar{z}_0, \dots, \bar{z}_{2m+1}]$  is an orientation-reversing diffeomorphism, and hence  $\mathbb{C}\mathbb{P}^{2m+1}$  and  $\overline{\mathbb{C}\mathbb{P}^{2m+1}}$  are diffeomorphic as oriented manifolds. Therefore, we see that  $N_k$  is diffeomorphic, as an oriented manifold, to the blowup of  $\mathbb{C}\mathbb{P}^{2m+1}$  at  $k-1$  points.

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