

CO781 / QIC 890:

Theory of Quantum Communication

Topic 2, part 3

The asymptotic equipartition theorem,  
Shannon entropy and classical data compression  
von Neumann entropy, Quantum data compression,

Entanglement concentration and dilution

Entropy of entanglement

Entanglement spread

Embezzlement of entanglement

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## References:

Preskill Sections 10.1.1, 10.3, 10.4

Entanglement concentration & dilution,  
entropy of entanglement

arXiv: quant-ph

Bennett, Bernstein, Popescu, Schumacher, 9511030

Lo, Popescu, 9902045

Lower bound for dilution & entanglement spread

Harrow, Lo, 0204096, Hayden, Winter, 0204092

Harrow, 0909.1557

## Embezzlement:

Hayden, van Dam 0201041

Leung, Toner, Watrous, 0804.4118

Consider the task:

For a given state  $|\psi\rangle \in \mathbb{C}^d \otimes \mathbb{C}^d$

(known to all of Richard, Alice, Bob)

Richard prepares  $|\psi\rangle^{\otimes n}$  on  $R_1A_1, R_2A_2, \dots, R_nA_n$

gives  $A_1 A_2 \dots A_n$  to Alice

Alice encodes  $A_1 \dots A_n$  into  $nr$  qubits

Alice sends those  $nr$  qubits to Bob

Bob decodes those  $nr$  qubits, output  $B_1 B_2 \dots B_n$ .

Requires final state on  $R_1 B_1 R_2 B_2 \dots R_n B_n \approx |\psi\rangle^{\otimes n}$

How to min  $r$  ?

Let  $\rho = \text{tr}_R |\psi\rangle\langle\psi|$ . TTS works, achieves the rate  $r = S(\rho)$ .

defined using  $\rho, n, \delta = r - S(\rho), \epsilon \dots$

Detail: use Schmidt decomposition  $|\psi\rangle = \sum_r \sqrt{p_r} |f_r\rangle_R |e_r\rangle_A$

$\rho = \sum_r p_r |e_r\rangle\langle e_r|$  is a spectral decomposition.

Let  $\epsilon > 0$ ,  $\delta = r - s(\rho) > 0$ ,  $T_{n,\delta}$  defined as before (on  $V$ ).

$$\Pi_S = \sum_{r^n \in T_{n,\delta}} |e_{r^n}\rangle\langle e_{r^n}|$$

Apply TTS to  $A_1 A_2 \dots A_n$ , output (on  $R_1 B_1 \dots R_n B_n$ ) =

$$\rho_{\text{out}} = \underbrace{I \otimes \Pi_S}_{\text{on } R^n A^n} \underbrace{\left( |\psi\rangle\langle\psi| \right)^{\otimes n}}_{R^n A^n} \underbrace{I \otimes \Pi_S}_{R^n A^n} + \text{tr} \left( I \otimes (I - \Pi_S) \left( |\psi\rangle\langle\psi| \right)^{\otimes n} \right) \cdot \text{ERR}$$

$$|\psi\rangle^{\otimes n} = \sum_{r^n} \sqrt{p_{r^n}} |f_{r^n}\rangle_{R^n} |e_{r^n}\rangle_{A^n}$$

$$(I \otimes \Pi_S) |\psi\rangle^{\otimes n} = \sum_{r^n \in T_{n,\delta}} \sqrt{p_{r^n}} |f_{r^n}\rangle_{R^n} |e_{r^n}\rangle_{A^n}$$

$$\langle\psi|^{\otimes n} (I \otimes \Pi_S) |\psi\rangle^{\otimes n} = \sum_{r^n \in T_{n,\delta}} p_{r^n} \geq 1 - \epsilon$$

$$\text{tr} \left( I \otimes (I - \Pi_S) \left( |\psi\rangle\langle\psi| \right)^{\otimes n} \right) \leq \epsilon$$

$$\begin{aligned}
& \left\| (\lvert \psi \rangle \langle \psi \rvert)^{\otimes n} - \rho_{\text{out}} \right\|_1 \\
&= \left\| (\lvert \psi \rangle \langle \psi \rvert)^{\otimes n} - \mathbb{I} \otimes \pi_S (\lvert \psi \rangle \langle \psi \rvert)^{\otimes n} \mathbb{I} \otimes \pi_S - \text{tr}(\mathbb{I} \otimes (\mathbb{I} - \pi_S) (\lvert \psi \rangle \langle \psi \rvert)^{\otimes n}) \cdot \text{ERR} \right\|_1 \\
&\leq \left\| (\lvert \psi \rangle \langle \psi \rvert)^{\otimes n} - \mathbb{I} \otimes \pi_S (\lvert \psi \rangle \langle \psi \rvert)^{\otimes n} \mathbb{I} \otimes \pi_S \right\|_1 + \left\| \text{tr}(\mathbb{I} \otimes (\mathbb{I} - \pi_S) (\lvert \psi \rangle \langle \psi \rvert)^{\otimes n}) \cdot \text{ERR} \right\|_1 \\
&\leq 2 \sqrt{1 - (\langle \psi \rvert^{\otimes n} (\mathbb{I} \otimes \pi_S) \lvert \psi \rangle^{\otimes n})^2} + \epsilon \leq 2\sqrt{2}\sqrt{\epsilon} + \epsilon.
\end{aligned}$$

↑  
notes on the detail last lecture ...

## Entanglement dilution and concentration (BBPS9511030)

Will see later why classical comm cannot create entanglement

LOCC is the class of operations consisting of unlimited  
Local Operations and Classical Communication

Two related questions in bipartite LOCC:

- (a) How many ebits are needed to create approx of  $|\Psi\rangle^{\otimes n}$  ?
- (b) How many ebits can be approx extracted from  $|\Psi\rangle^{\otimes n}$  ?

Approx in trace distance (for pure states, same as fidelity up to square root).

Task (a) is called entanglement dilution

Task (b) is called entanglement concentration

Answer for (a) upper bounds answer for (b).

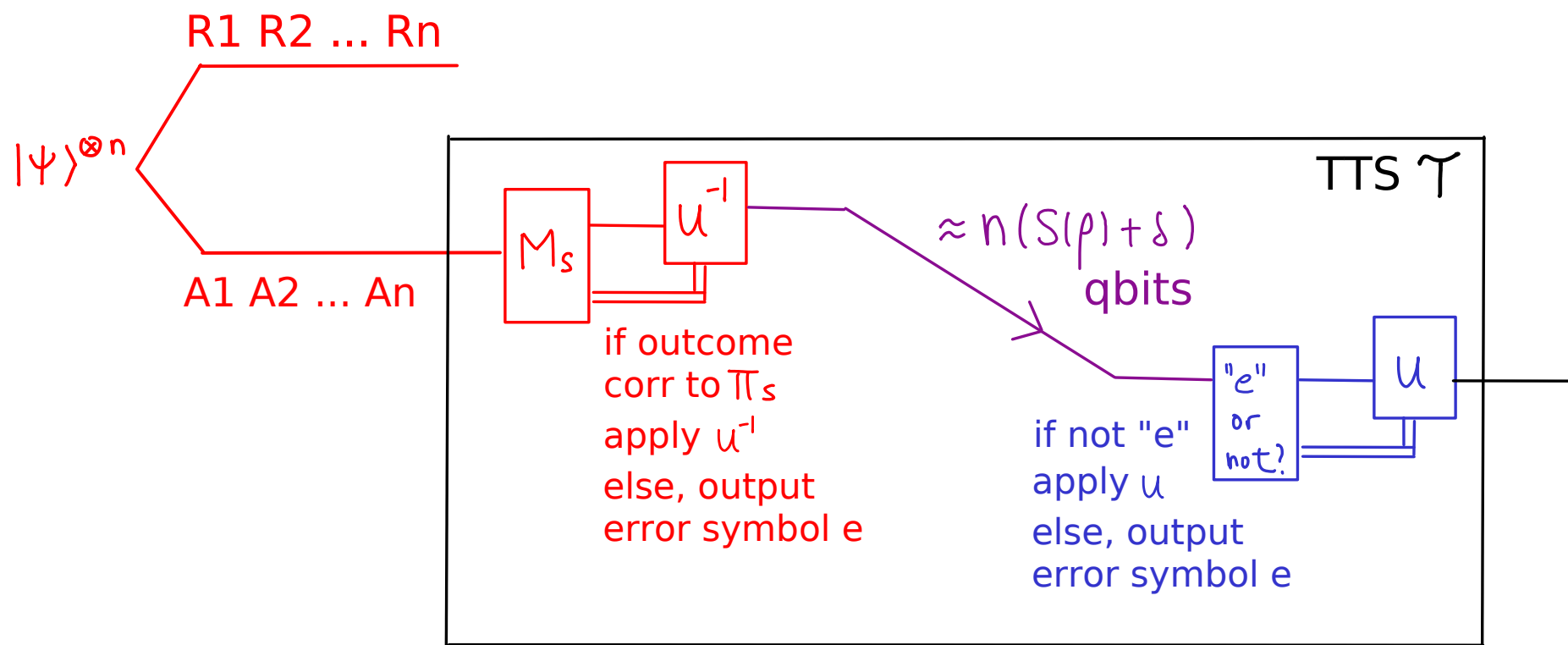
Nice surprise: both ans are  $nS(\rho) + o(n)$ ,  $\rho = \text{tr}_B |\Psi\rangle\langle\Psi|$ .

(Note both Alice and Bob know what states they are transforming to and from.)

## For entanglement dilution:

Simplest method: use the protocol to distribute  $|\psi\rangle^{\otimes n}$  but obtain the qubits needed by teleportation.

In detail: here Alice is also Richard. She prepares  $|\psi\rangle^{\otimes n}$ .

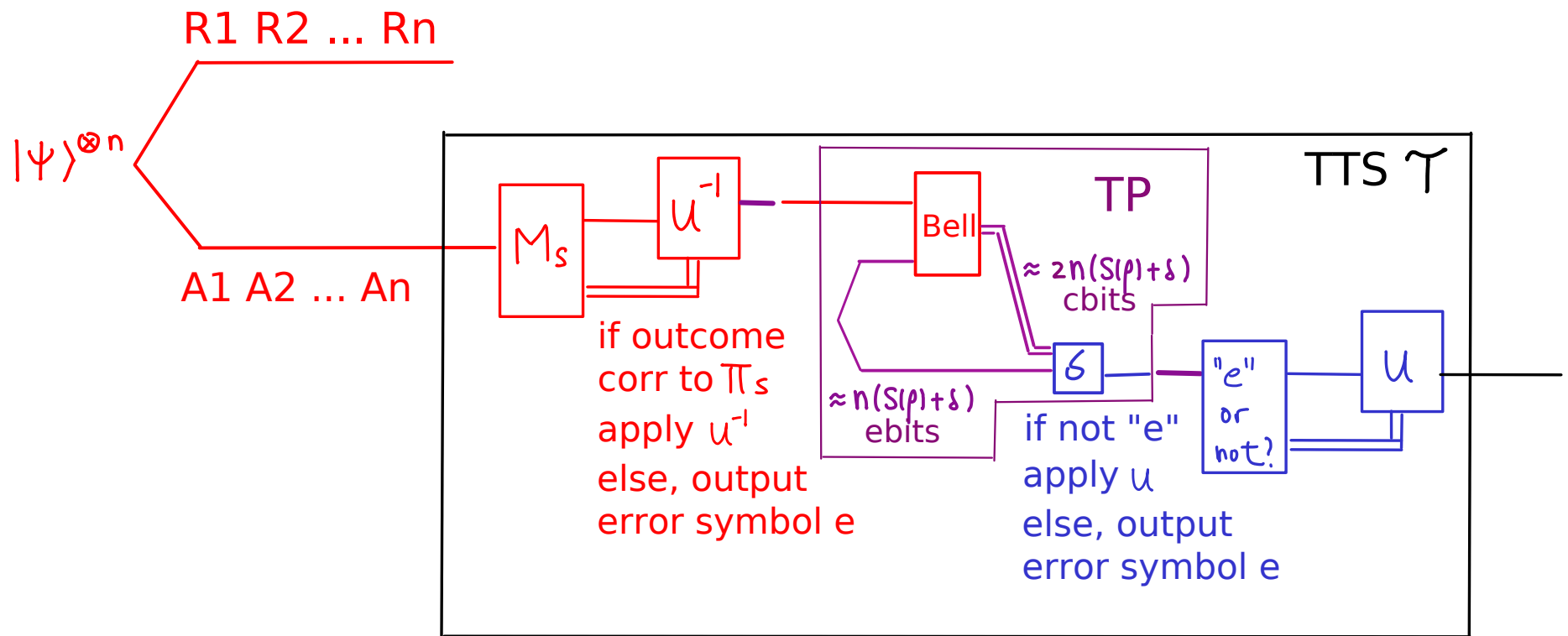


We already proved the above outputs a state close to  $|\psi\rangle^{\otimes n}$ .  
Now replace qbit by TP (doesn't change the output).

## For entanglement dilution:

Simplest method:  $n(S(\rho) + \delta)$  ebits +  $2n(S(\rho) + \delta)$  cbits  $\geq |\Psi\rangle^{\otimes n}$ .  
free in LOCC

In detail: here Alice is also Richard. She prepares  $|\Psi\rangle^{\otimes n}$ .



We already proved the above outputs a state close to  $|\Psi\rangle^{\otimes n}$ .  
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# How to reduce the classical comm cost for dilution?

Main idea (Lo Popescu) (detail in A2):

up to local unitaries

$$|\psi\rangle^{\otimes n} \approx a \frac{1}{\sqrt{2}} (|00\rangle + |11\rangle) \otimes n S(p) - O(\sqrt{n}) \otimes |\mu\rangle + b |\nu\rangle$$

Method:

goes away  
when  $b |\nu\rangle$   
ignored

use  
 $n S(p) - O(\sqrt{n})$   
ebits

teleport using  
 $O(\sqrt{n})$  ebits  
 $O(\sqrt{n})$  cbits

ignore

$O(\sqrt{n})$  qubit state  
small

## For entanglement concentration:

$$\text{Let } |\psi\rangle = \sum_{\nu} \sqrt{p_{\nu}} |e_{\nu}\rangle_A |f_{\nu}\rangle_B$$

Local unitaries are free under LOCC, so, Alice and Bob first convert the Schmidt basis to computational basis for each copy. They now have  $|\phi\rangle^{\otimes n}$  where

$$|\phi\rangle = \sum_{\nu} \sqrt{p_{\nu}} |\nu\rangle_A |\nu\rangle_B$$

$$|\phi\rangle^{\otimes n} = \sum_{\nu^n} \sqrt{p_{\nu^n}} |\nu^n\rangle_{A^n} |\nu^n\rangle_{B^n}$$

Tempting idea: each of Alice & Bob applies meas w/ POVM  $\{\pi_s, I - \pi_s\}$ . With prob  $\geq 1 - \epsilon$ , obtain state

$$|\widehat{\phi}^n\rangle \propto \sum_{\nu^n \in T_{n,\delta}} \sqrt{p_{\nu^n}} |\nu^n\rangle_{A^n} |\nu^n\rangle_{B^n}$$

While  $2^{-n(S(p)+\delta)} \leq p_{\nu^n} \leq 2^{-n(S(p)-\delta)}$  ( $p_{\nu^n}$ 's roughly equal)

A2:  $|\widehat{\phi}^n\rangle$  is NOT close to being maximally entangled :(

For entanglement concentration:

Alice and Bob will make a much finer measurement.

Def: Let rv  $X$  has sample space  $\{1, 2, \dots, m\}$ .

Let  $\mathcal{X}^n = \mathcal{X}_1 \mathcal{X}_2 \dots \mathcal{X}_n$  be the outcome for  $n$  iid draws.

Suppose  $t_k = \# \mathcal{X}_i$ 's equal to  $k$ .

Then,  $\mathcal{X}^n$  is in the **type class**  $(t_1, t_2, \dots, t_m)$ .

**$m$ -tuples of non-neg integers summing to  $n$**

e.g.,  $n$  coin tosses ( $m=2$ ), all outcomes with  $k$  1's and  $(n-k)$  2's are in the type class  $(k, n-k)$ .

e.g., 20 throws of a dice ( $n=20, m=6$ ).

Outcome 44326511564314622246 has

$t_1 = 3$ ,  $t_2 = 4$ ,  $t_3 = 2$ ,  $t_4 = 5$ ,  $t_5 = 2$ ,  $t_6 = 4$

so, outcome is in the type class  $(3, 4, 2, 5, 2, 4)$ .

Outcomes in the same type class are exactly equiprobable.

For entanglement concentration:

Alice and Bob each measures the type class  
(Alice measures  $A_1 A_2 \dots A_n$ , Bob  $B_1 B_2 \dots B_n$ ).

They always get the same outcome.

Conditioned on each outcome, their postmeas state is maximally entangled, with Schmidt rank depending on which type class.

In A2, you will show that the expected # ebits is

$$n S(\rho) - o(n)$$

for  $|\psi\rangle \in \mathbb{C}^2 \otimes \mathbb{C}^2$ .

Idea holds for general  $d$ , if  $n$  large enough.

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(1) Operational meaning:

for bipartite pure entangled state in many copies

- LOCC conversion is approx reversible
- single "currency" (ebit) of entanglement

This gives meaning to the quantity  $S(\text{tr}_B |\Psi\rangle\langle\Psi|)$  as "the amount of entanglement" in the state  $|\Psi\rangle$ .

(2) Even better, above holds even if CC is charged

- concentration requires no CC
- dilution requires  $\Theta(\sqrt{n})$  cbits

achievability: Lo-Popescu, necessity: Harrow-Lo-Hayden-Winter

(3) For bipartite pure state single copy, (1)-(2) don't hold.

LOCC conversion: Lo-Popescu 9703038

Nielsen 9811053 (majorization)

(4) For bipartite mixed state, many copies

Task (a) has no name (?). # ebits per copy required:

entanglement of formation Bennett DiVincenzo Smolin Wootters 96

regularized to "entanglement cost" Hayden (M) Horodecki

and the two can be different

Terhal 0008134

Shor 0305035, Hastings 0809.3972

restricting to vanishing CC, # ebits per copy is called

the entanglement of purification.

Terhal Horodecki Leung

DiVincenzo 0202044

Task (b) is called entanglement distillation

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# ebits extracted per copy: distillable entanglement

Mix state has "noise" ... to be removed by distillation.

Distillation (with 1- or 2-way CC) is mathematically equivalent to noisy channel coding for sending quantum data through noisy quantum channels (+crypto apps)!

Also, distillable entanglement can be strictly smaller than entanglement cost for some state (e.g., "bound entangled states" have 0 distillable entanglement but positive entanglement cost). M, P, R Horodecki 9801069

So, no single entanglement measure for mixed state, and LOCC conversion can be irreversible.

(5) For 3 or more parties, pure state, large # of copies  
no comparable conversion theory  
many types of incomparable entanglement

Bennett Popescu Rohrlich Smolin Thapliyal 9908073



## Entanglement spread:

Def: for bipartite state  $|\Psi\rangle_{AB}$ ,  $\rho = \text{tr}_B |\Psi\rangle\langle\Psi|$ ,

\* its entanglement spread is defined as

$$\Delta(|\Psi\rangle) = \log(\text{rank}(\rho)) - \log \frac{1}{\|\rho\|_\infty}$$

\* its  $\epsilon$  perturbed entanglement spread is defined as

$$\Delta_\epsilon(|\Psi\rangle) = \min_{\substack{P: \text{projectors} \\ \text{tr}(\rho P) \geq 1-\epsilon}} \Delta(P \otimes I |\Psi\rangle)$$

e.g.,  $\Delta = 0$  iff all nonzero Schmidt coeffs are equal.

The transformation:  $|\phi\rangle \rightarrow |\tilde{\psi}\rangle$

s.t. fidelity of  $|\psi\rangle, |\tilde{\psi}\rangle \geq 1 - \epsilon$

requires  $C \geq \Delta_{(4\epsilon)^{\frac{1}{2}}}(|\psi\rangle) - \Delta_0(|\phi\rangle) + 2 \log(1 - (4\epsilon)^{\frac{1}{2}})$  cbits.

i.e., increase in spread must be "paid for" by classical communication.

For entanglement dilution:

$$|\bar{\phi}\rangle^{\otimes nr} \longrightarrow |\psi\rangle^{\otimes n}$$

↑  
ebits

0 spread

↑  
 $\mathcal{O}(n)$  spread

$\Theta(\sqrt{n}) \in$  perturbed spread

**lower bound of cbits needed**