

Least-Squares Finite Element Methods for Nonlinear Hyperbolic Conservation Laws

3 April 2003

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Nonlinear Hyperbolic Conservation Laws

$$\begin{aligned} \nabla_{t,x} \cdot (u, f(u)) &= 0 & \text{in } \Omega \\ u &= g & \text{on } \Gamma_{in} \end{aligned}$$

$$\nabla = (\partial_t, \partial_x) \quad \nabla^\perp = (-\partial_x, \partial_t)$$

inviscid Burgers equation:

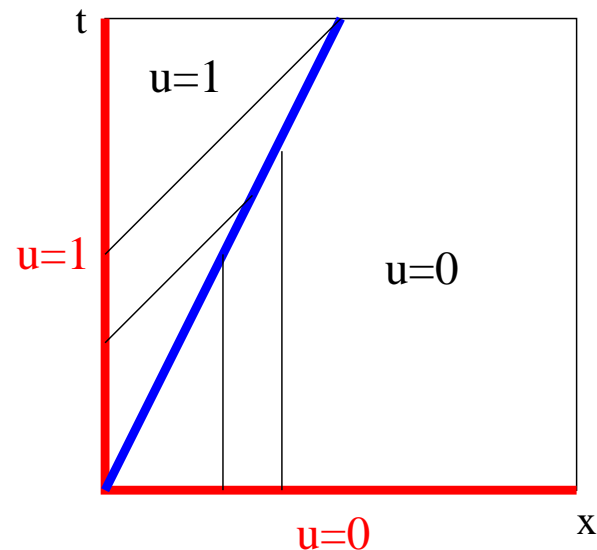
$$f(u) = u^2/2$$

Rankine-Hugoniot relation:

$$[(u, f(u))] \cdot \vec{n} = 0$$

$$\Rightarrow (u, f(u)) \in H(\text{div})$$

$$H(\text{div}) = \{(u, v) \in (L_2)^2 \mid \|\nabla \cdot (u, v)\|_0^2 < \infty\}$$



Outline

- Least-squares finite element method (LSFEM)
- (1) LSFEM for the Burgers equation
- (2) $H(\text{div})$ -conforming LSFEM
- (3) Dual $H(\text{div})$ -conforming LSFEM
- FEM convergence theory
- Numerical conservation
- Adaptivity
- Conclusions



Least-Squares Finite Element Method

- $Lu = f$

- define the functional

$$\mathcal{F}(u; f) = \|Lu - f\|_0^2 = \langle Lu - f, Lu - f \rangle$$

⇒ minimization:

$$u_*^h = \arg \min_{u^h \in \mathcal{U}^h} \|Lu^h - f\|_0^2 = \arg \min \mathcal{F}(u^h; f)$$

- condition for stationary point:

$$\frac{\partial \mathcal{F}(u_*^h + \alpha v^h; f)}{\partial \alpha} \Big|_{\alpha=0} = 0 \quad \forall v^h \in \mathcal{U}^h$$

if L is linear: $\mathcal{F}(u_*^h + \alpha v^h; f) =$

$$\langle Lu_*^h - f, Lu_*^h - f \rangle + 2\alpha \langle Lu_*^h - f, Lv^h \rangle + \alpha^2 \langle v^h, v^h \rangle$$

⇒ weak form:

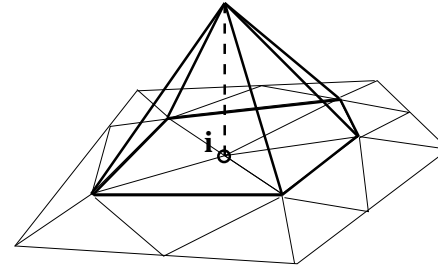
$$\text{find } u_*^h \in \mathcal{U}^h, \text{ s.t. } \langle Lu_*^h, Lv^h \rangle = \langle f, Lv^h \rangle \quad \forall v^h \in \mathcal{U}^h$$



Finite Element Discretization

- approximate $u \in \mathcal{U}$ by $u^h \in \mathcal{U}^h$

$$u^h(t, x) = \sum_{i=1}^n u_i \phi_i(t, x)$$



- algebraic system from weak form:

$$\langle Lu^h, L\phi_j \rangle = \langle f, L\phi_j \rangle \quad \forall \phi_j$$

equation j :
$$\sum_{i=1}^n u_i \langle L\phi_i, L\phi_j \rangle = \langle f, L\phi_j \rangle$$

(n equations in n unknowns)



(1) LSFEM for the Burgers Equation

- Gauss-Newton: first linearize the equation, then put the linearized equation into the LS functional, then minimize the LS functional
- Linearization:

$$F(u) := \nabla \cdot (u, u^2/2) = 0$$

$$\text{Newton: } F(u) = 0 \Rightarrow F(u_0) + dF|_{u_0}(u - u_0) = 0$$

The Fréchet derivative $dF|_{u_0}(v)$ at u_0 in a direction v is

$$dF|_{u_0}(v) = \lim_{\varepsilon \rightarrow 0} \frac{F(u_0 + \varepsilon v) - F(u_0)}{\varepsilon}$$

$$dF|_{u_0}(v) = \nabla \cdot ((1, u_0) v)$$

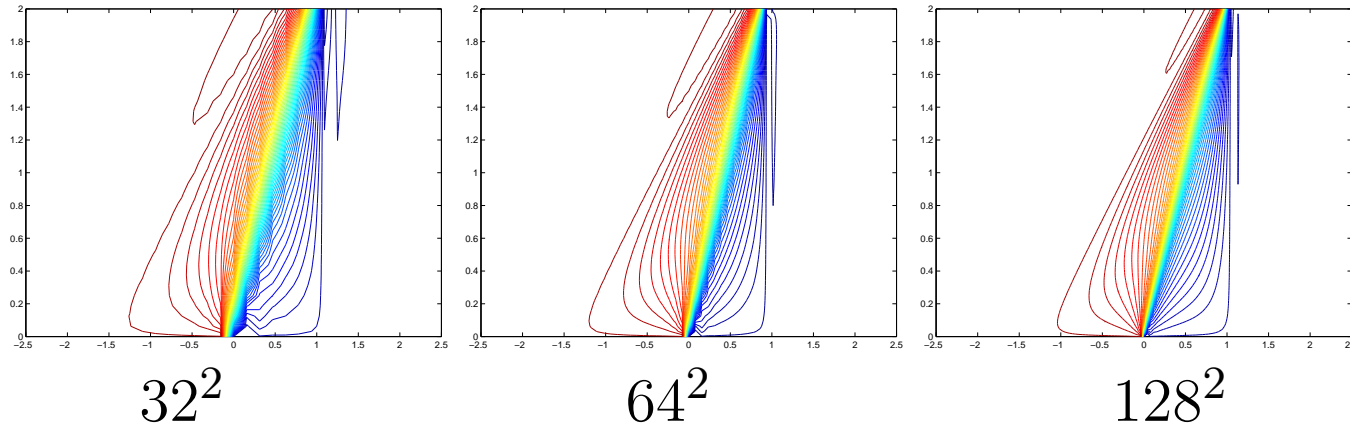
$$\Rightarrow \nabla \cdot (u_0, u_0^2/2) + \nabla \cdot ((1, u_0) (u - u_0)) = 0$$



Numerical Results

$$u_*^h = \arg \min_{u^h \in \mathcal{U}^h} \|\nabla \cdot (u, u^2/2)\|^2$$

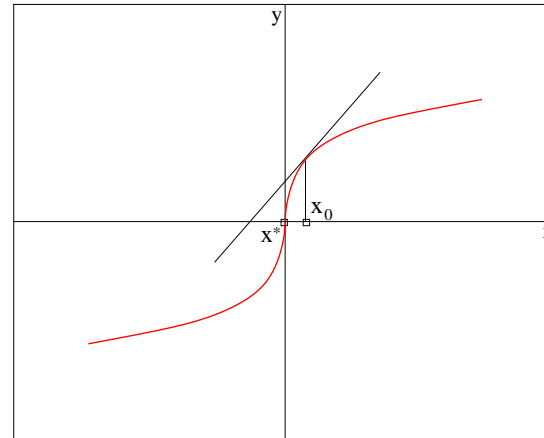
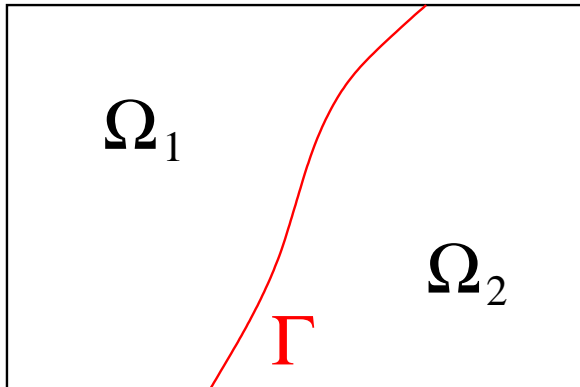
$u^h \in \mathcal{U}^h \subset H^1$, bilinear elements on quadrilaterals



- right shock speed, no oscillations
- on each grid, newton process converges
- BUT: for decreasing h , functional does not go to zero
- this means: for decreasing h , convergence to an incorrect solution!!! (L^*L has a spurious stationary point)
- WHY NO CONVERGENCE??



Unbounded Fréchet Derivative



$$\nabla \cdot \vec{w} = \nabla \cdot \vec{w}_{1,2}^{(p)} + [\vec{w}]_1^2 \cdot \vec{n} \delta_\Gamma$$

$$dF|_{u_0}(v) = \nabla \cdot ((1, u_0) v) = \nabla \cdot ((1, u_0) v)_{1,2}^{(p)} + [(1, u_0) v]_1^2 \cdot \vec{n} \delta_\Gamma$$

- when u_0 is discontinuous (e.g. u_0 is the exact solution), for almost all functions v the Fréchet derivative is unbounded!! (recall: $[(u_0, u_0^2/2)]_1^2 \cdot \vec{n} = 0$)

- Newton with $f'(x_*) = \infty$ may have **empty basin of attraction** (e.g. $f(x) = |x|^{1/3} \Rightarrow x_1 = -2x_0$)

\Rightarrow this may explain why LSFEM fails to converge (convergence to a spurious stationary point)



(2) $H(\text{div})$ -conforming Reformulation

$$\begin{aligned} \nabla \cdot (u, u^2/2) &= 0 & \Omega \\ u &= g & \Gamma_{in} \end{aligned} \quad \Rightarrow$$

$$\begin{aligned} \nabla \cdot \vec{w} &= 0 & \Omega \\ \vec{w} - (u, u^2/2) &= 0 \\ \vec{n} \cdot (\vec{w} - (g, g^2/2)) &= 0 & \Gamma_{in} \end{aligned}$$

$$\vec{F}(\vec{w}, u) = 0 \quad \Rightarrow \quad \vec{F}(\vec{w}_0, u_0) + \vec{F}'|_{\vec{w}_0, u_0}(\vec{w} - \vec{w}_0, u - u_0) = 0$$

$$\vec{F}'|_{\vec{w}_0, u_0}(\vec{w} - \vec{w}_0, u - u_0) = \begin{bmatrix} \nabla \cdot & 0 \\ I & -1 \\ & -u_0 \end{bmatrix} \cdot \begin{bmatrix} \vec{w} - \vec{w}_0 \\ u - u_0 \end{bmatrix}$$

\Rightarrow Fréchet derivative is **bounded!**
(we choose $\vec{w} \in H(\text{div})$)



(3) Dual $H(\text{div})$ -conforming Reformulation

$$\begin{aligned} \nabla \cdot (u, u^2/2) &= 0 & \Omega \\ u &= g & \Gamma_{in} \end{aligned} \quad \Rightarrow \quad \begin{aligned} \nabla^\perp p - (u, u^2/2) &= 0 & \Omega \\ \vec{n} \cdot (\nabla^\perp p - (g, g^2/2)) &= 0 & \Gamma_{in} \end{aligned}$$

$$\vec{F}(p, u) = 0 \quad \Rightarrow \quad \vec{F}(p_0, u_0) + \vec{F}'|_{p_0, u_0}(p - p_0, u - u_0) = 0$$

$$\vec{F}'|_{p_0, u_0}(p - p_0, u - u_0) = \begin{bmatrix} \nabla^\perp & -1 \\ & -u_0 \end{bmatrix} \cdot \begin{bmatrix} p - p_0 \\ u - u_0 \end{bmatrix}$$

\Rightarrow Fréchet derivative is **bounded!** (we choose $p \in H(\text{curl})$)

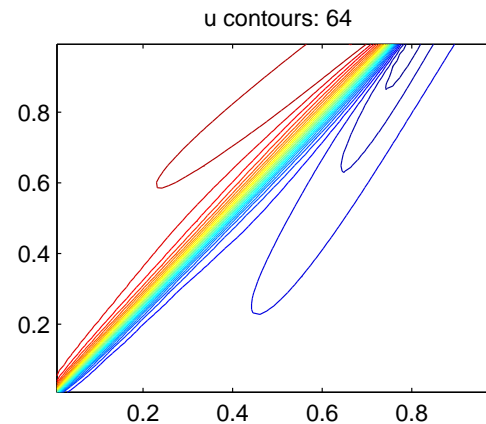
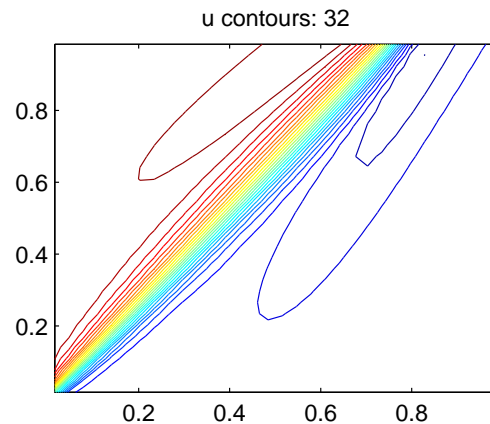
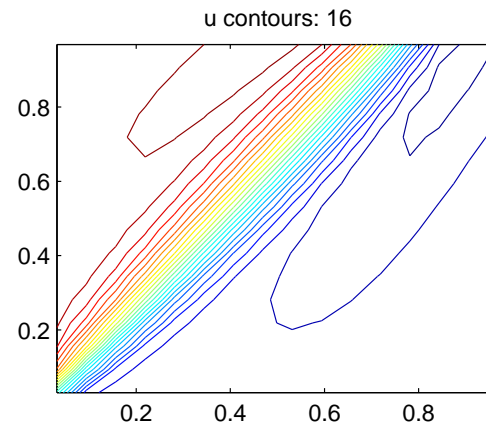
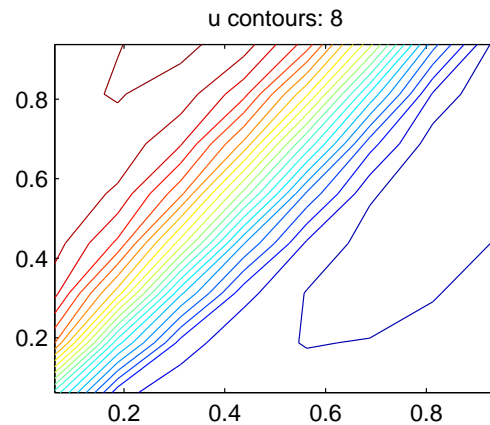
- p is **De Rham-dual** of \vec{w}
 - move up one space to the left in De Rham-diagram of differential forms
 - similar to potential formulations



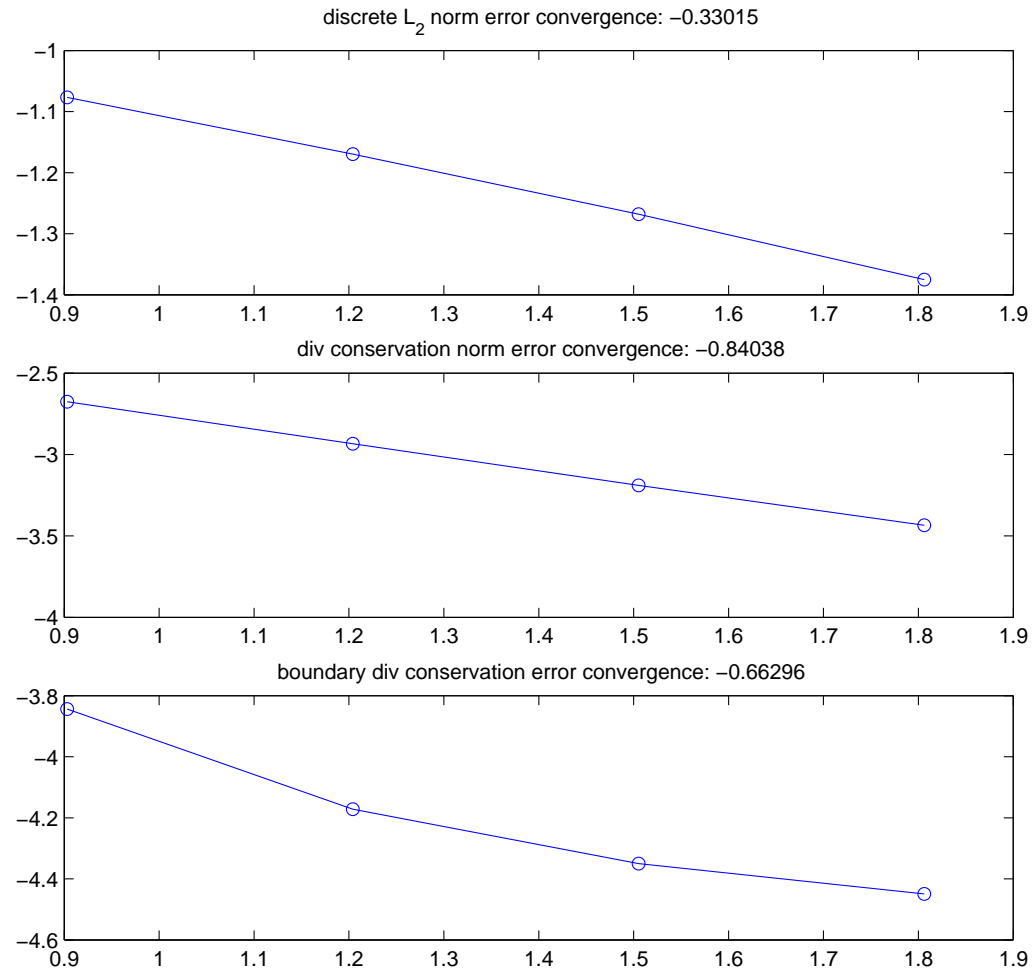
Results for $H(\text{div})$ -conforming LSFEM

$$(\vec{w}_*^h, u_*^h) = \arg \min_{\vec{w}^h \in \mathcal{W}_g^h, u^h \in \mathcal{U}^h} \|\nabla \cdot \vec{w}\|^2 + \|\vec{w} - (u, u^2/2)\|^2$$

\vec{w}^h Raviart-Thomas elements $\subset H(\text{div})$, u^h bilinear elements on quadrilaterals, strong boundary conditions



L2 convergence rate to exact solution

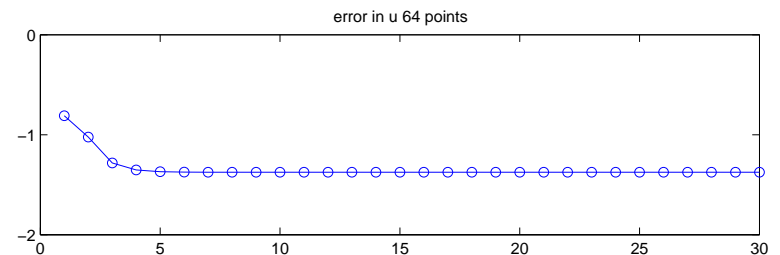
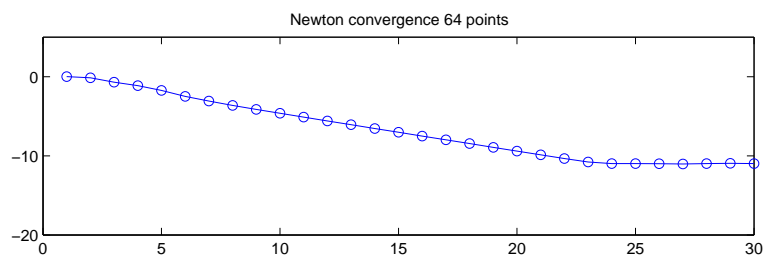
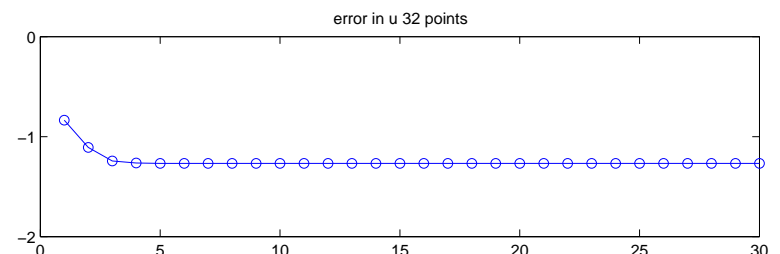
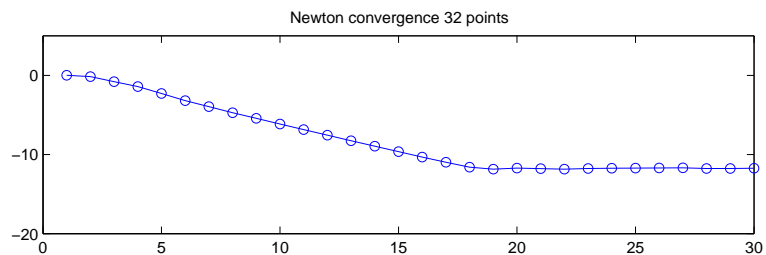
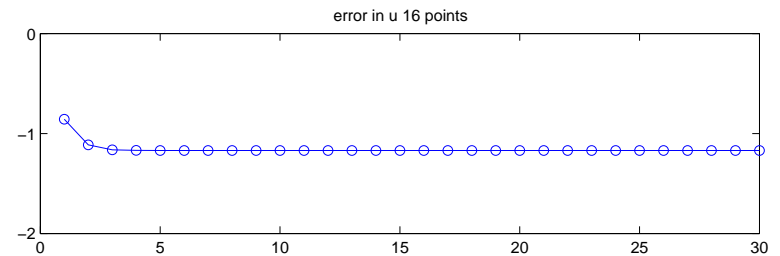
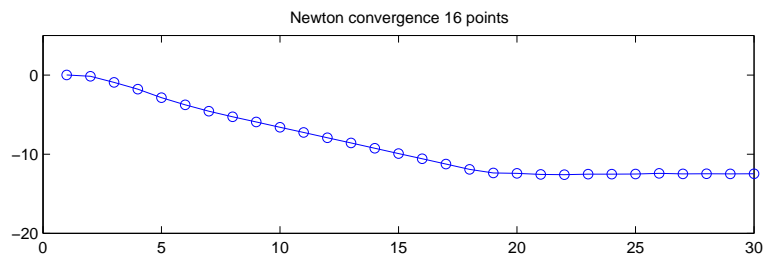
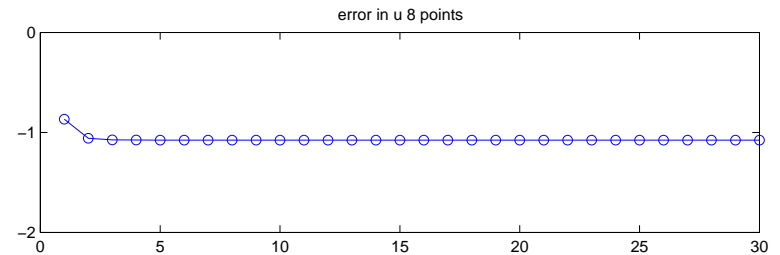
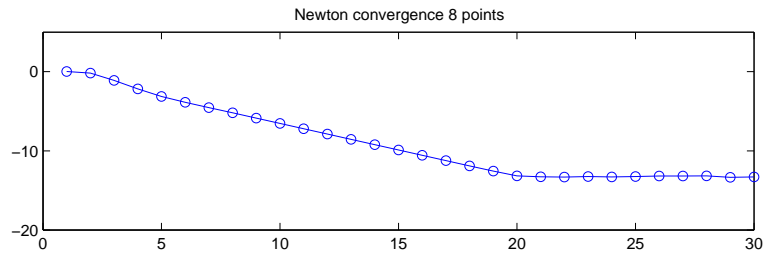


$$\|u^h - u_{exact}\|_0^2 = O(h^{0.66}) \quad (\text{optimal} : O(h^1))$$

$$\|\nabla \cdot \vec{w}\|_0^2 = O(h^{1.68})$$



Gauss-Newton Convergence

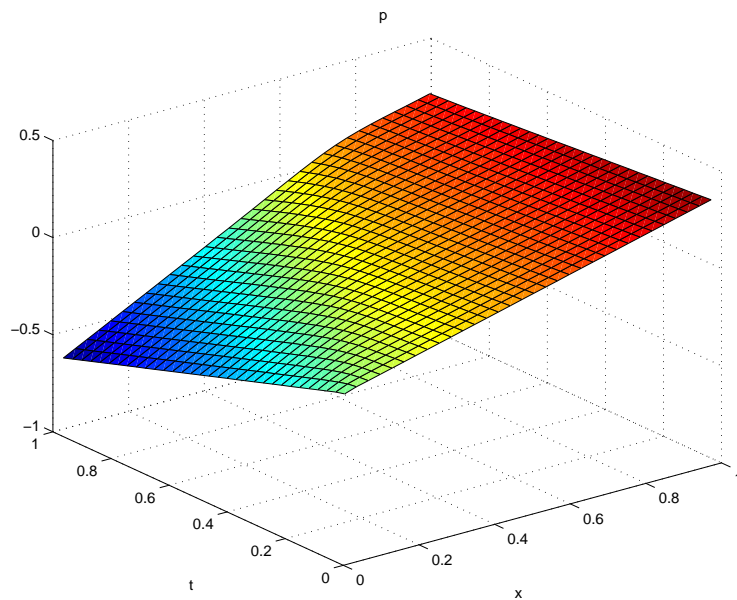


Results for dual $H(\text{div})$ -conforming LSFEM

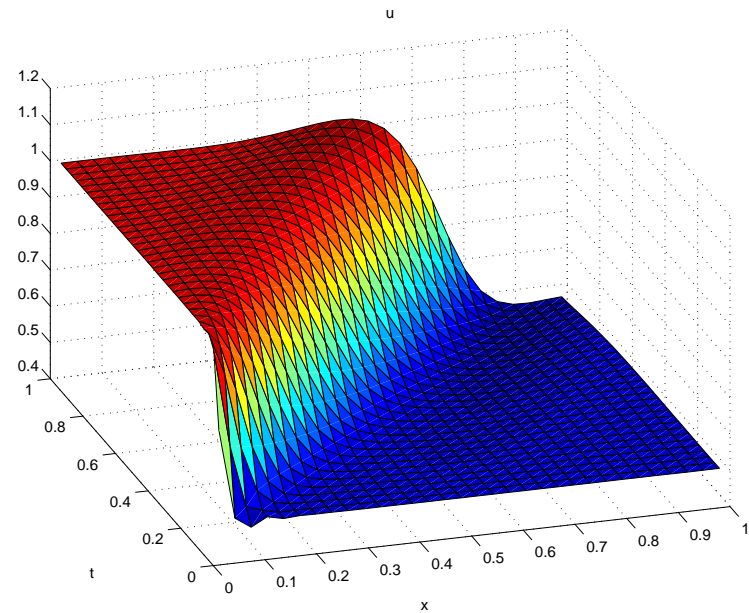
$$(p_*^h, u_*^h) =$$

$$\arg \min_{p^h \in \mathcal{P}^h, u^h \in \mathcal{U}^h} \|\nabla^\perp p - (u, u^2/2)\|^2 + \|\vec{n} \cdot (\nabla^\perp p - (g, g^2/2))\|_B^2$$

p^h, u^h bilinear elements on quadrilaterals, weak boundary conditions



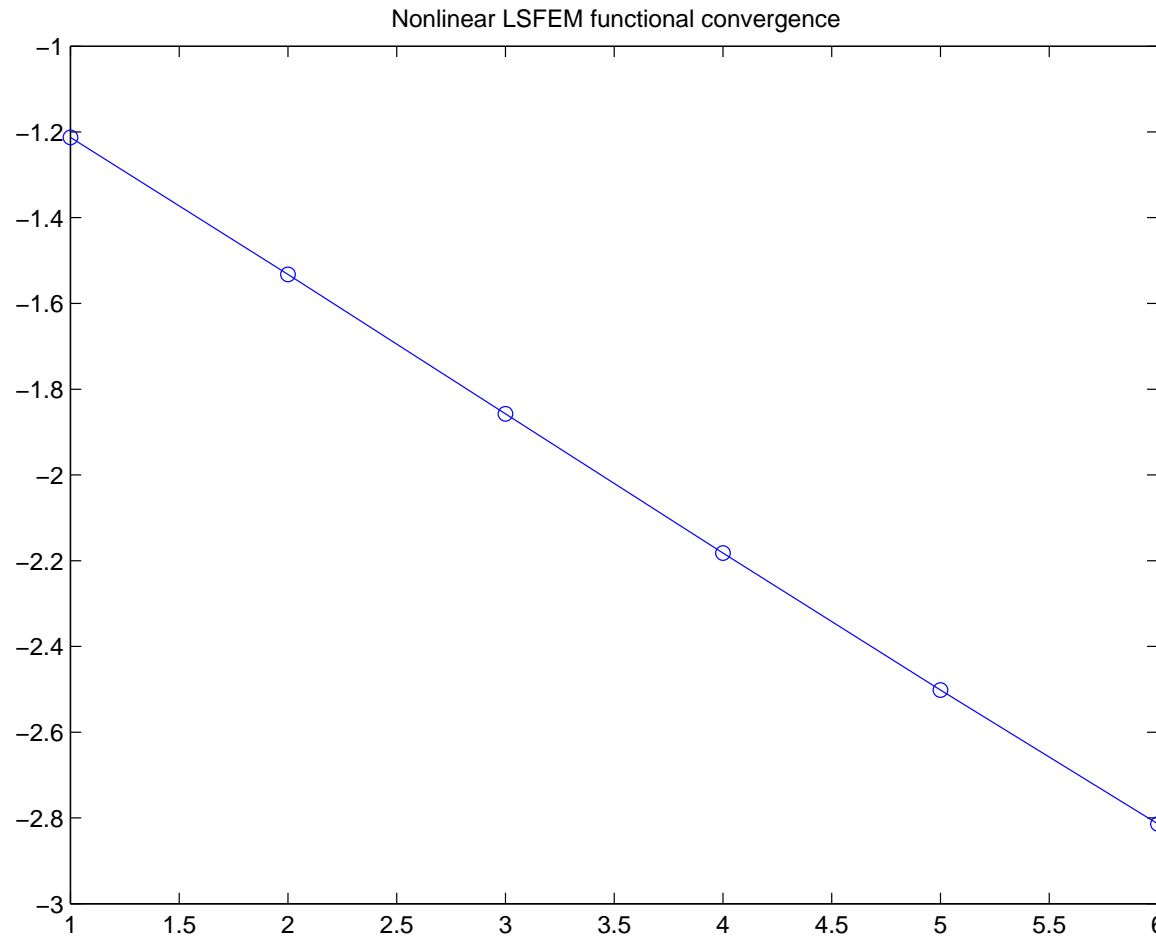
p



u



Nonlinear functional convergence



$$\|\nabla^\perp p - (u, f(u))\|^2 + \|\vec{n} \cdot (\nabla^\perp p - (g, f(g)))\|_B^2 = O(h^{1.1})$$



Convergence of $H(\text{div})$ -conforming LSFEM

- is the LS functional uniformly coercive?

$$\exists c \text{ s.t. } \|\nabla^\perp p - (u, f(u))\|^2 \geq c (\|\nabla^\perp p\|^2 + \|u\|^2)$$

NO: there are **high-frequency error modes** for which

$$c = O((1/n)^2)$$

- **compatible spaces**: if we choose both p^h and u^h **bilinear**, we find (1D heuristics and numerical evidence) that $c = O(h)$ (only for the low-frequency modes, high-frequency error is filtered out)

remark: if we choose u^h piecewise constant, high-frequency error contaminates solution



Convergence of $H(\text{div})$ -conforming LSFEM

- we observe that $\|\nabla^\perp p^h - (u^h, f(u^h))\|^2$ converges faster than $\|\nabla^\perp p^h\|^2 + \|u^h\|^2$

⇒ conjecture:

$$\|\nabla^\perp p^h - (u^h, f(u^h))\|^2 \geq c h (\|\nabla^\perp p^h\|^2 + \|u^h\|^2)$$

with

$$\|\nabla^\perp p^h - (u^h, f(u^h))\|^2 = O(h^{1+\beta})$$

$$\|\nabla^\perp p^h\|^2 + \|u^h\|^2 = O(h^\alpha)$$

- to be investigated further



Numerical Conservation

THEOREM. Lax-Wendroff (1960).

'conservative' finite difference formula:

$$\frac{u_i^{h,n+1} - u_i^{h,n}}{\Delta t} + \frac{\bar{f}_{i+1/2}^{h,n} - \bar{f}_{i-1/2}^{h,n}}{\Delta x} = 0,$$

scheme converges \implies convergence to weak solution

\implies **exact discrete conservation** ($\nabla \cdot (u, f(u))_d \equiv 0$) is a *sufficient* condition for convergence to a weak solution, but is often *erroneously* considered as *necessary*

- popular FEM for hyperbolic conservation laws (e.g. Discontinuous Galerkin) are discretely conservative in the Lax-Wendroff sense



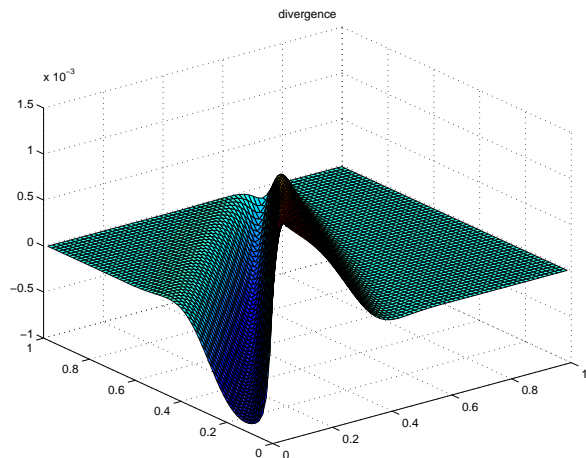
Numerical Conservation

THEOREM. Conservation theorem for $H(\text{div})$ -conforming LSFEM.

$H(\text{div})$ -conforming LSFEM converges \implies convergence to weak solution

THEOREM. Conservation theorem for dual $H(\text{div})$ -conforming LSFEM.

dual $H(\text{div})$ -conforming LSFEM converges \implies convergence to weak solution



- $H(\text{div})$ -conforming LSFEM does not impose strict discrete numerical conservation, but converges to weak solution! \implies discrete numerical conservation is not necessary

- dual $H(\text{div})$ -conforming LSFEM has stronger, *pointwise* discrete conservation property: $\nabla \cdot (\nabla^\perp p^h) \equiv 0$



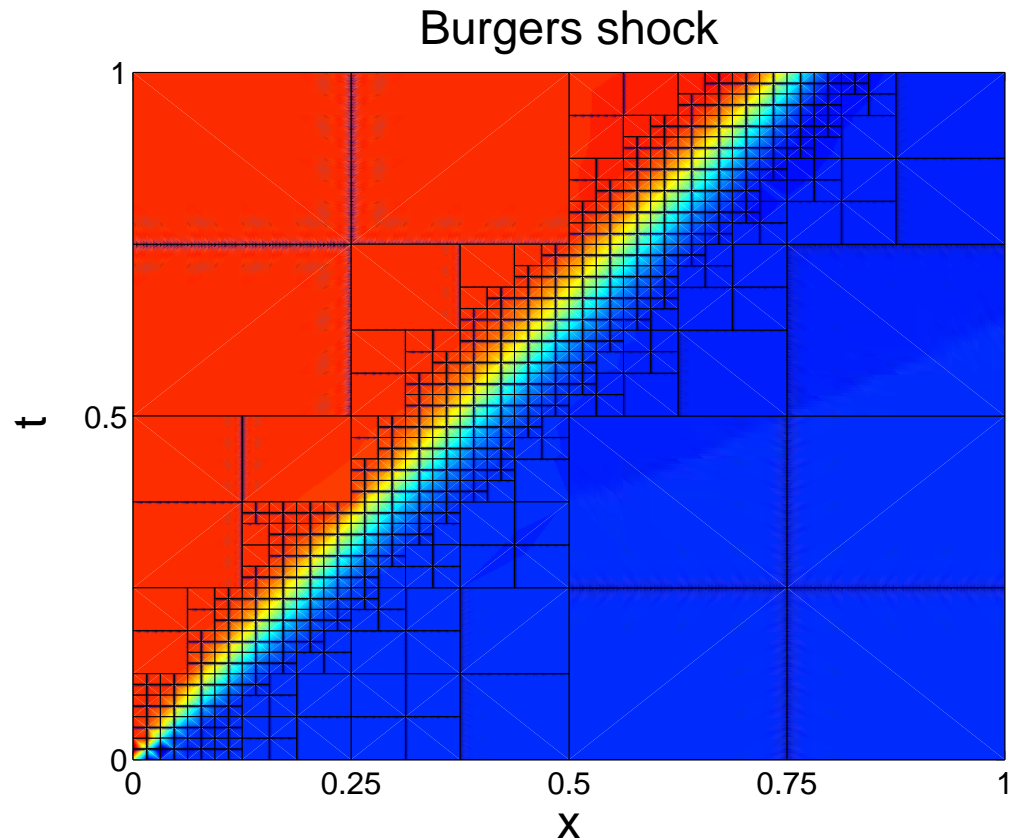
Error Estimator and Adaptive Refinement

$$\begin{aligned}\mathcal{F}(u^h; f) &= \|Lu^h - f\|_0^2 \\ &= \|Lu^h - Lu_{exact}\|_0^2 \\ &= \|L(u^h - u_{exact})\|_0^2 \\ &= \|Le^h\|_0^2\end{aligned}$$

- functional value gives sharp local a posteriori error estimator
- use error estimator for adaptive refinement in space–time



Error Estimator and Adaptive Refinement



- sharp fronts at shocks
- with optimal $O(n)$ solver, work per grid point is bounded
- promises to be competitive with other methods (also explicit timemarching)



Conclusions

we have developed two classes of $H(\text{div})$ -conforming LSFEM for hyperbolic conservation laws

- disadvantages

- extra variables are introduced (\vec{w} or p)
- high diffusion of LSFEM at shocks

- advantages of LSFEM

- optimal solution within finite element space
- SPD linear systems (iterative methods, AMG)
- error estimator (efficient adaptive refinement)
- conservation (either weaker or stronger than Lax-Wendroff discrete conservation)
- no spurious oscillations at discontinuities (without need to add numerical diffusion)
- linear higher order schemes



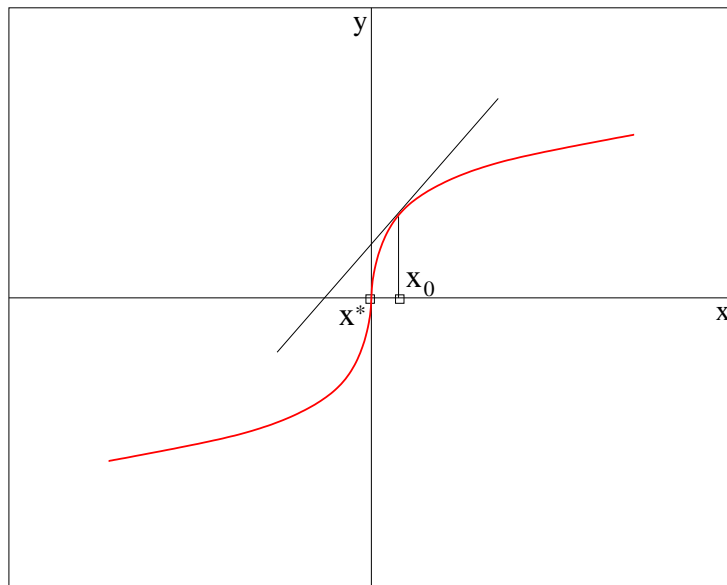
Conclusions

- advantages of $H(\text{div})$ -conforming reformulation
 - linear differential operator
 - bounded Fréchet derivative \Rightarrow Newton converges
 - nullspace of operator can be represented exactly
 - differential boundary conditions help AMG (dual)
 - regularity of the solution ($\vec{w} \in H(\text{div})$) is made explicit, also at the discrete level using Raviart-Thomas elements
- FE convergence theory remains to be worked out further
- promising initial AMG results, to be developed further
- methods can be extended to multiple spatial dimensions, and to systems of equations



Newton with Unbounded Derivative

$$f(x) = 0 \Rightarrow f(x_0) + f'(x_0)(x - x_0) = 0 \Rightarrow x = x_0 - \frac{f(x_0)}{f'(x_0)}$$



$$f(x) = x^\alpha \quad x > 0$$

$$= -|x|^\alpha \quad x \leq 0$$

$$\alpha \in (0, 1)$$

$$f'(x) = \alpha x^{\alpha-1} \quad x > 0$$

$$f'(0) = \infty$$

$$x = x_0 - \frac{x_0^\alpha}{\alpha x_0^{\alpha-1}}$$

$$x = (1 - 1/\alpha) x_0$$

divergence $\forall \alpha < 1/2$, e.g. for $\alpha = 1/3$, $x = -2 x_0$

\Rightarrow if $f'(x^*) = \infty$, basin of attraction can be empty!



Algebraic Multigrid

- use approach investigated for linear problems for $H(\text{div})$ -conforming LSFEM

- e.g., dual LSFEM: $L = \nabla^\perp p - \vec{b}u$

$$\langle \nabla^\perp p^h, \nabla^\perp q^h \rangle + \langle -\vec{b}u^h, \nabla^\perp q^h \rangle = 0 \quad \forall q^h \in \mathcal{P}^h$$

$$\langle \nabla^\perp p^h, -\vec{b}s^h \rangle + \langle -\vec{b}u^h, -\vec{b}s^h \rangle = 0 \quad \forall s^h \in \mathcal{U}^h$$

$$A = \begin{bmatrix} A_{pp} & A_{pu} \\ A_{up} & A_{uu} \end{bmatrix}$$

- A is symmetric positive definite
- standard AMG efficient for A_{pp}
- A_{uu} mass matrix, but A_{up} strong off-diagonal coupling
- some promising initial results, but work in progress



Numerical Conservation

$$\partial_t(u_i^h S_i) + \sum_{j \in \partial\Omega_i} \bar{f}_j^h \cdot \vec{n}_j l_j = 0.$$

$$\partial_t\left(\sum_{i \in \Omega_s} u_i^h S_i\right) + \sum_{j \in \partial\Omega_s} \bar{f}_j^h \cdot \vec{n}_j l_j = 0.$$

